

Riemann-Roch formula (not the most general form, we'll discuss sheaf, sheaf cohomology, ... to develop necessary tools later).

- $p_1, \dots, p_d \in X$ ,  $D := \{p_1, \dots, p_d\} \subseteq X$ .
- $H^0(D) := \{ \text{mero. fcn. on } X \text{ w/ at worst simple poles at } D \}$ .
- $H^0(K-D) := \{ \text{holo. 1-forms on } X \text{ that vanish at each } p_i \}$ .

Then:  $h^0(D) - h^0(K-D) = d - g + 1$ .

- recall that residue is not well-defined for mero fcn, but is well-defined for mero. 1-forms. Given  $f \in H^0(D)$ , at each  $p_i \in D$ , consider.

$$\begin{array}{ccc} T_{p_i}^* X & \longrightarrow & \mathbb{C} \\ \omega & \longmapsto & \text{res}_{p_i}(f\omega) \end{array}$$

which can be considered as an element of  $T_{p_i}^* X$ .

$$\Rightarrow \text{Res}: H^0(D) \longrightarrow \bigoplus_{i=1}^d T_{p_i}^* X$$

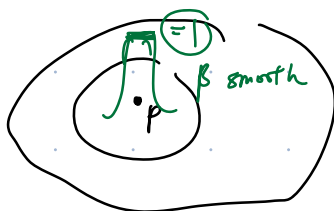
- $\text{Ker}(\text{Res}) = \{ \text{holo. fcn. on } X \} = \{ \text{constant fcn.} \}$ .
- The image of Res is the local residue data at  $D$  that can be realized by a global mero. fcn. on  $X$ .

Recall that, this is closely related to  $H^1(X)$ , which gives the obstruction of such local data to glue globally.

Let's define a map  $T_p X \longrightarrow H^{0,1}(X)$  as follows:

- Let  $z$  be a local coord. at  $p$ , an element of  $T_p X$  is of the form  $a \frac{\partial}{\partial z}$ .

- Consider a bump fn,  $= 1$  near  $p$ ,  $= 0$  outside of a small disc.



- Then  $A := \bar{\partial} \left( \beta \cdot \frac{a}{z} \right) = \underbrace{(\bar{\partial} \beta)}_{\substack{\text{"near } p \\ 0}} \cdot \frac{a}{z}$  defines a global  $(0,1)$ -form.

\*  $\frac{a}{z}$  is indep. of the choice of coord. !!

- $[A] \in H^{0,1}(X)$  is indep. of the choice of  $\beta$ .

This defines a map:  $T_p X \longrightarrow H^{0,1}(X)$ .

$$a \frac{\partial}{\partial z} \longmapsto \left[ \bar{\partial} \beta \cdot \frac{a}{z} \right]$$

We have an exact seq.:

$$0 \rightarrow \mathbb{C} \xrightarrow{\text{const.}} H^0(D) \xrightarrow{\text{Res}} \bigoplus_{i=1}^d T_{p_i} X \xrightarrow{IA} H^{0,1}(X).$$

(the given local residue data in  $\bigoplus T_{p_i} X$  can be repr. by global mero. fcn.  $\Leftrightarrow$  their image in  $H^{0,1}(X)$  sum to 0.)

The dual map:  $A^T: H^{0,1}(X)^* \xrightarrow{\text{HIS}} \bigoplus_{i=1}^d T_{p_i}^* X$   
 $H^{1,0}(X)$

Linear alg. Fact:

$$\begin{aligned} \dim \ker A - \dim \ker A^T &= d - \dim H^{0,1}(X) = d - g. \\ \parallel \\ \dim \operatorname{Im}(Res) \\ \parallel \\ h^0(D) - 1. \end{aligned}$$

Claim:  $A^T: H^{1,0}(X) \longrightarrow \bigoplus_{i=1}^d T_{p_i}^* X$  is  $(2\pi i) \cdot ev.$   
 where  $ev: H^{1,0}(X) \longrightarrow \bigoplus T_{p_i}^* X$   
 $\omega \longmapsto (\omega|_{p_i})$

Assuming the Claim, then

$$\begin{aligned} \dim \ker A^T &= \dim \{ \text{holo. 1-forms vanishing at } D \} \\ &= h^0(K-D). \end{aligned}$$

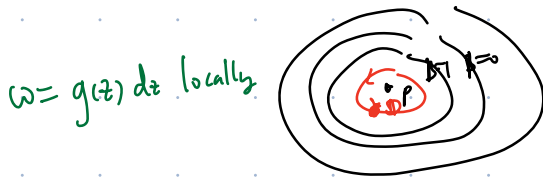
Which proves Riemann-Roch.

pf of Claim:  $T_p X \longrightarrow H^{0,1}(X)$ ,  $H^{0,1}(X)^* \longrightarrow T_p^* X$   
 $a \frac{\partial}{\partial \bar{z}} \longmapsto \bar{\partial}(\beta \cdot \frac{a}{\bar{z}})$   
 $\omega \in H^{1,0} \longmapsto (T_p X \longrightarrow \mathbb{C})$   
 $\alpha \frac{\partial}{\partial \bar{z}} \longmapsto \bar{\partial}(\beta \frac{a}{\bar{z}})$   
 $([\gamma] \in H^{0,1}) \longmapsto \int \omega \wedge \gamma$   
 $\int \omega \wedge \bar{\partial}(\beta \cdot \frac{a}{\bar{z}})$

$$\int \omega \wedge \bar{\partial}(\beta \frac{a}{\bar{z}}) = a \int \omega \wedge \bar{\partial}(\beta) \frac{1}{\bar{z}}$$

$$= a \int_{X \setminus D} \omega \wedge (\bar{\partial} \beta) \frac{1}{\bar{z}} = a \int_{X \setminus D} \bar{\partial}(\omega \wedge \frac{\beta}{\bar{z}}) = a \int_{X \setminus D} d(\omega \wedge \frac{\beta}{\bar{z}})$$

holo 1-form on  $X \setminus D$



$$\stackrel{\text{Stokes}}{=} a \int_{\gamma} \omega \wedge \frac{\beta}{\bar{z}} = a \int_{\gamma} \frac{\omega}{\bar{z}} = 2\pi i \cdot a \cdot g(p)$$

$$\begin{array}{ccc} \text{ev: } H^{1,0}(X) & \longrightarrow & T_p^* X \\ \omega & \longmapsto & (T_p X \rightarrow \mathbb{C}) \\ \text{"} & & a \frac{\partial}{\partial z} \mapsto a g_{0,1} \\ g_{(1,1)d?} & & \end{array} \quad \square$$

Last time:

- $H^{1,0}(X) \cong \overline{H^{0,1}(X)}$ ;  $\alpha \mapsto [\bar{\alpha}]$ .
- $H^{1,0}(X) \cong H^{0,1}(X)^*$
- $H_{dR}^1(X) \cong H^{1,0}(X) \oplus H^{0,1}(X)$ ;  $H_{dR}^2(X) \cong H^{1,1}(X)$ .

Refs: More generally, for a cpt. Kähler mfd  $X$ , we have the Hodge decomposition:

$$H_{dR}^k(X, \mathbb{C}) \cong \bigoplus_{p+q=k} H^{p,q}(X).$$

Our proofs are based on the following thm:

Thm:  $X$ : cpt. R.S.  $\gamma \in A^2(X)_{\mathbb{C}}$ .

"  $\exists f \in A^0(X)_{\mathbb{C}}$  st.  $\Delta f = \gamma$ "  $\iff$  " $\int_X \gamma = 0$ ".

Hodge thm:

- every class in  $H_{dR}^k(M)$  can be uniquely represented by a harmonic  $k$ -form (denoted  $H^k(M)$ )
- There is an orthogonal decomposition:  

$$A^k(M) = H^k(M) \oplus \Delta(A^k(M))$$

(we'll discuss later!)

In the thm.,  $(\Rightarrow)$  is easy:

$$\int_X \Delta f = 2i \int_X \bar{\partial} \partial f = 2i \int_X d(\partial f) = 0.$$

$(\Leftarrow)$  is the most crucial part of the thm. It can be directly obtained from Hodge thm:

- $H^2(X) \cong H_{\text{dR}}^2(X) \cong H^2(X)$  is  $1\text{-dim}^{\mathbb{R}}$ , spanned by the volume form of  $X$ .

We can repr it by a harmonic 2-form  $\omega$ . we have  $\int_X \omega \neq 0$ .

- $\mathcal{G} \in A^2(X) = H^2(X) \oplus \Delta(A^2(X))$

i.e.  $\exists c \in \mathbb{C}, f \in A^0(X)$  s.t.  $\mathcal{G} = c \cdot \omega + \Delta f$ .

$$\Rightarrow 0 = \int_X \mathcal{G} = \underbrace{c \int_X \omega}_\neq 0 + \underbrace{\int_X \Delta f}_0$$

$$\Rightarrow c = 0. \quad \square$$

Remark: Further generalization: "Elliptic Diff<sup>1,2</sup> operator" on cpt mfd:

$L$  w/ adjoint  $L^*$ :  $\langle Lf, g \rangle = \langle f, L^*g \rangle$ .

$L$  elliptic  $\Rightarrow A^k(M) = \ker L^* \oplus L A^k(M)$   
 $\uparrow$   
 finite dim <sup>$\mathbb{R}$</sup>

$\Delta^* = \Delta \Rightarrow A^k(M) = H^k(M) \oplus \Delta A^k(M)$

- It suffices to prove the real version of the main thm.

$$\int_{\mathcal{D}} \rho = 0 \Rightarrow \exists f \text{ s.t. } \Delta f = \rho$$

Def:  $f, g \in C^\infty(X)$

- Dirichlet inner product:

$$\begin{aligned} \langle f, g \rangle_{\mathcal{D}} &:= 2i \int_X \partial f \wedge \bar{\partial} g \\ &= - \int_X \left( \frac{\partial f}{\partial x} \frac{\partial g}{\partial x} + \frac{\partial f}{\partial y} \frac{\partial g}{\partial y} \right) dx \wedge dy \\ &= \int_X g \Delta f = \int_X f \Delta g \end{aligned}$$

- Dirichlet norm:  $\|f\|_{\mathcal{D}}^2 := \langle f, f \rangle_{\mathcal{D}} = \int_X f \Delta f = \int_X |df|^2$ .

- They actually define on the quotient:  $C^\infty(X)/\mathbb{R}$   $\leftarrow$  constant functions.

Fact:  $(C^\infty(X)/\mathbb{R}, \langle \cdot, \cdot \rangle_{\mathcal{D}})$  is a pre-Hilbert space. (not complete)

Let  $g \in A^2(X)_{\mathbb{R}}$ .  $\forall \phi, \psi \in C^\infty(X)$ ,

$$\int_X \psi (g - \Delta \phi) = \int_X \psi g - \langle \phi, \psi \rangle_{\mathcal{D}}$$

$$\| \Delta \phi = g \| \Leftrightarrow \| \int_X \psi (g - \Delta \phi) = 0 \quad \forall \psi \|$$

$$\Leftrightarrow \| \int_X \psi g = \langle \phi, \psi \rangle_{\mathcal{D}} \quad \forall \psi \|$$

Main Thm can be formulated as: If  $\int_X g = 0$ ,

(then  $\hat{f}: C^\infty(X)/\mathbb{R} \longrightarrow \mathbb{R}$  is well-defined,  
 $\psi \longmapsto \int_X \psi g$ )

$$\exists \phi \in C^\infty(X) \text{ s.t. } \hat{f}(\psi) = \langle \psi, \phi \rangle_D \quad \forall \psi.$$

Riesz repr thm: Let  $H$  be a Hilbert space,

$\varphi: H \longrightarrow \mathbb{R}$  be a bounded linear map ( $|\varphi(x)| \leq C \|x\|$ ).

Then:  $\exists z \in H$  s.t.  $\varphi(x) = \langle x, z \rangle \quad \forall x \in H.$

• Any pre-Hilbert space has a formal completion: a point in the completed Hilbert space is an equiv. class of Cauchy seq.  $\{\psi_i\}$ , where  $\{\psi_i\} \sim \{\psi'_i\}$  if  $\|\psi_i - \psi'_i\| \rightarrow 0$ .

• Need 1: For  $\int g = 0$ , the functional  $\hat{f}: C^\infty(X)/\mathbb{R} \longrightarrow \mathbb{R}$  is bounded,  
 $\psi \longmapsto \int \psi g$

i.e.  $\exists C > 0$  s.t.  $|\int \psi g| \leq C \|\psi\|_D \quad \forall \psi \in C^\infty(X).$

• Such functional extends naturally to a functional on  $H$ : (completion of  $C^\infty(X)/\mathbb{R}$ ).

• By Riesz repr thm,  $\exists \phi \in H$  s.t.  $\hat{f}(\psi) = \langle \psi, \phi \rangle_D \quad \forall \psi.$  "weak sol<sup>n</sup>"

• Need 2: For  $\int g = 0$ , such weak sol<sup>n</sup> is actually an element of  $C^\infty(X)/\mathbb{R}$ . ( $\subseteq H$ ). (regularity of sol<sup>n</sup> of ell. PDE)

Boundedness of  $\hat{g}$ :  $(|\int_X \gamma \beta| \leq C \|\psi\|_D = \int_X |d\psi|^2)$

(Need: Bound the  $L^2$ -norm of  $\psi$  by the  $L^2$ -norm of  $d\psi$ .)

" $L^2$ -Poincaré ineq."  $\Omega$ -bdd. convex domain in  $\mathbb{R}^2$ ,  
 $\psi$ -sm. fun. on an open set containing  $\overline{\Omega}$ .

$$\int_{\Omega} |\psi|^2 d\mu \leq \frac{d^2}{4\pi^2} \int_{\Omega} |d\psi|^2 d\mu, \quad d = \text{diam}(\Omega).$$

Local case: Suppose  $\beta$  is supported on a single coord. chart.  $\simeq \Omega$ .

Then  $\beta$  can be identified with a smooth fun. on  $\Omega$ .

$$|\hat{g}(\psi)| = \left| \int_X \beta \psi \right| = \left| \int_{\Omega} \beta \psi d\mu \right| \leq \|\beta\|_{L^2(\Omega)} \|\psi\|_{L^2(\Omega)} \quad (\text{Cauchy-Schwarz})$$

$$\leq \underbrace{\|\beta\|_{L^2(\Omega)} \cdot \frac{\text{diam}(\Omega)}{\pi}}_C \cdot \underbrace{\|\psi\|_D}_{\leftarrow \|\psi\|_{L^2(\Omega)} \text{ by def.}^2 \text{ of Dirichlet norm.}}$$

General  $\beta$  w/  $\int_X \beta = 0$ :

• Since  $H_{\mathbb{R}}^2(X) \xrightarrow{\sim} \mathbb{R}$ ,  $[\beta] = 0$  in  $H_{\text{dR}}^2$ , i.e.  $\beta = d\theta$  for some  $\theta \in A^1(X)$ .

• Let  $\{U_{\alpha}\}$  be a (finite) cover of  $X$  by coord. charts, w/ partition of unity  $\{\chi_{\alpha}\}$  supported on each  $U_{\alpha}$ .

• Then:  $\beta_{\alpha} := d(\chi_{\alpha}\theta)$  is supp. on  $U_{\alpha}$ , w/  $\int_X \beta_{\alpha} = 0$ .

$$\Rightarrow |\hat{\beta}_{\alpha}(\psi)| = \left| \int_X \beta_{\alpha} \psi \right| \leq C_{\alpha} \|\psi\|_D, \quad \forall \psi \in C^{\infty}(X).$$

$$\begin{aligned} \hat{g}(\psi) &= \int_X \beta \psi = \int_X (d\theta) \psi = \int_X d(\sum \chi_{\alpha} \theta) \psi = \sum \int_X \beta_{\alpha} \psi \\ &= \sum_{\alpha} \hat{\beta}_{\alpha}(\psi) \Rightarrow |\hat{g}(\psi)| \leq (\sum C_{\alpha}) \|\psi\|_D. \quad \square \end{aligned}$$

finite sum.

The weak sol<sup>n</sup> is actually smooth.  $\hat{f}(\psi) = \langle \psi, "f" \rangle$

• By Riesz repr. thm, we know "f" exists in the completion of  $C^\infty(X)/\mathbb{R}$ ,

i.e.  $\exists$  seq.  $\{\phi_i\} \subseteq C^\infty(X)/\mathbb{R}$ . Cauchy w.r.t. the Dirichlet norm,

with:  $\langle \phi_i, \psi \rangle \longrightarrow \hat{f}(\psi)$  as  $i \rightarrow \infty$ .

Claim:  $\{\phi_i\}$  conv. to a function which is locally  $L^2$ .

• Note that there is a choice  $\phi_i + (\text{const.})$  involved, so we need to be a bit more careful.

• Let  $\Omega$  be a coord. chart. We choose suitable constants s.t.  $\int_{\Omega} \phi_i d\mu = 0$

•  $L^2$ -Poincaré ineq.  $\Rightarrow \|\phi_i - \phi_j\|_{L^2(\Omega)} \leq C(\Omega) \cdot \|\phi_i - \phi_j\|_{\mathcal{D}}$ .

• Since  $\{\phi_i\}$  is Cauchy in Dirichlet norm,  $\{\phi_i\}$  is a Cauchy seq. in  $L^2(\Omega)$ .

By completeness of  $L^2(\Omega)$ ,  $\{\phi_i\} \longrightarrow \phi \in L^2(\Omega)$ .

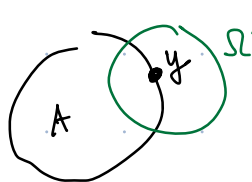
• We claim that the same seq.  $\{\phi_i\}$  (with the chosen constants), converges locally in  $L^2$  over all of  $X$ . This uses the standard open/closed argument.

$A := \{x \in X \mid \exists \text{ local chart } \Omega \text{ of } x \text{ s.t. } \phi_i \rightarrow \phi \in L^2(\Omega)\}$ .

By previous discussion,  $A \neq \emptyset$ . It's clear that  $A$  is open.

• It remains to show that  $A$  is closed. Suppose  $y \in \bar{A} \setminus A$ .

We know that  $\exists y \in \Omega'$  and  $\{c_i\}$  s.t.  $\phi_i - c_i \longrightarrow \tilde{\phi} \in L^2(\Omega')$ .


 For  $x \in A \cap \Omega'$ , in a small nbhd of  $x$ ,  
 both  $\{\phi_i\}$  &  $\{\phi_i - c_i\}$  conv. in  $L^2$

$\Rightarrow c_i \rightarrow c$  for some  $c$ .

$\Rightarrow \{\phi_i\} \rightarrow \tilde{\phi} + c$  in  $L^2(\Omega')$ , thus  $y \in A$ .  $\square$

We have so far:  $\exists \phi$ : locally  $L^2$ -fcn. on  $X$  s.t.

$$\hat{g}(\psi) = \langle \psi, \phi \rangle \quad \forall \psi \in \mathcal{E}^\infty(X)$$

$$\stackrel{=}{=} \int_X \psi \cdot g \quad \stackrel{=}{=} \int_X (\Delta \psi) \cdot \phi.$$

i.e. a locally  $L^2$ -sol<sup>n</sup> of " $\Delta \phi = g$ ".

Want:  $\phi \in \mathcal{E}^\infty(X)$ .

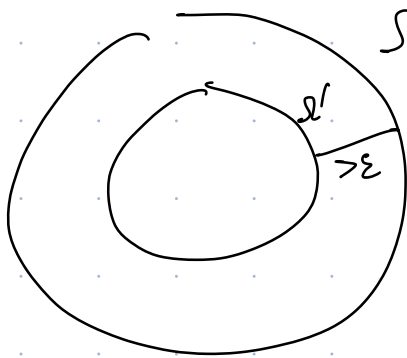
This is a purely local statement of  $\phi$ . so we can just consider coord. charts  $\Omega$  and sm. fcn's supp. on  $\Omega$ .

Thm:  $\Omega$ : bdd. open set in  $\mathbb{C}$ ,  $g$ : sm. 2-form on  $\Omega$ .

Suppose  $\phi \in L^2(\Omega)$  has:

$$\int_\Omega (\Delta \chi) \phi = \int_\Omega \chi \cdot g \quad \forall \chi \in \mathcal{E}_c^\infty(\Omega).$$

Then  $\phi$  is smooth.



It suffices to show the smoothness of  $\phi$  in  $\Omega'$ .

Let  $g'$  be a smooth fcn. st.  $g' = g$  in  $\Omega'$ ,  
 $\text{Supp}(g')$  is cpt and in  $\Omega$ .

$$\Rightarrow \int_{\Omega'} (\Delta \chi) \cdot \phi = \int_{\Omega'} \chi \cdot g' \quad \forall \chi \in \mathcal{C}_c^\infty(\Omega')$$

Fact: Let  $K$  be the kernel fcn:  $K(x) := \frac{1}{2\pi} \log|x|$ .

(not defined at 0, but is locally integrable fcn. on  $\mathbb{R}^2$ ).

- $(K * -)$  is the "inverse" of  $\Delta(-)$ :
  - If  $\sigma \in \mathcal{C}_c^\infty(\mathbb{R}^2)$ , then  $K * (\Delta \sigma) = \sigma$ .
  - If  $f \in \mathcal{C}_c^\infty(\mathbb{R}^2)$ , then  $\Delta(K * f) = f$ .  
*also smooth.*

where  $(K * f)(x) = \int_{\mathbb{R}^2} K(y) f(x-y) d\mu_y$  is the convolution.

$$\Rightarrow \exists \psi' \in \mathcal{C}^\infty \text{ st. } \Delta \psi' = g' \quad (\psi' = K * g')$$

$$\Rightarrow \int_{\Omega'} (\Delta \chi) \phi = \int_{\Omega'} \chi \cdot \Delta \psi' = \int_{\Omega'} (\Delta \chi) \cdot \psi' \quad \forall \chi \in \mathcal{C}_c^\infty$$

It remains to show: If  $\int_{\Omega'} (\Delta \chi) \cdot \phi = 0 \quad \forall \chi \in \mathcal{C}_c^\infty$ ,

then  $\phi$  is smooth.

(i.e.  $\phi$  is a "weak sol<sup>n</sup>" of the heat eq<sup>n</sup>  $\Delta \phi = 0$ ).

- Using the mean-value property of harmonic fns, one can construct an explicit sm. fn.  $B \in \mathcal{C}_c^\infty(\mathbb{R}^2)$  supp. on  $|z| \leq \varepsilon$ . s.t.

" Suppose  $\psi \in \mathcal{C}^\infty(\mathbb{R}^2)$  and  $\text{Supp}(\Delta \psi) \subseteq \underset{\text{cpt}}{J} \subseteq \mathbb{C}$ .

Then  $B * \psi - \psi$  is supported in the  $\varepsilon$ -nbhd of  $J$ ."

- Observe that if  $\phi$  is a sm. sol<sup>n</sup> of  $\Delta \phi = 0$ , then we have

$$B * \phi = \phi \text{ in } \Omega'.$$

- $\forall L^2$ -fcn  $\phi$ ,  $B * \phi \in \mathcal{C}^\infty$ .

Therefore, proving  $\phi \in \mathcal{C}^\infty \iff$  proving  $B * \phi = \phi$  in  $\Omega'$ .

- It suffices to show:  $\forall \chi \in \mathcal{C}_c^\infty(\Omega')$ ,

$$\int_{\Omega'} \chi \cdot (B * \phi - \phi) = 0.$$

- $K * \chi$  is smooth, with  $\Delta(K * \chi) = \chi \rightarrow$  supp. in  $\Omega'$ .

$\Rightarrow \underline{B * K * \chi - K * \chi}$  supp. in  $\Omega$  (containing  $\varepsilon$ -nbhd of  $\Omega'$ )

!!  
h

$$\Delta h = B * \chi - \chi.$$

$$\begin{aligned}
 0 &= \int_{\Omega} (\Delta \eta) \phi = \int_{\Omega} (B^* \chi - \chi) \phi \\
 &= \int_{\Omega} \chi (B^* \phi - \phi). \quad \square
 \end{aligned}$$

Rmk: (on harmonic forms).

$M$ : cpt. Riemannian mfd  $\mapsto$  inner product on  $A^k(M)$

Heuristic: In a class  $[\alpha] \in H_{dR}^k(M)$ , the infimum of  $\|\cdot\|$  of a repr of  $[\alpha]$  is achieved by harmonic form:  $(\Delta \alpha = 0)$

Suppose  $\alpha$  achieves the min. of  $\|\alpha + d\beta\|$ .

$$\text{Then } \|\alpha + t d\beta\|^2 \geq \|\alpha\|^2$$

$$\|\alpha\|^2 + 2t \langle \alpha, d\beta \rangle + t^2 \langle d\beta, d\beta \rangle$$

holds for all  $t \in \mathbb{R}$ .  
 $\beta \in A^{k-1}$

$$\Rightarrow \langle \alpha, d\beta \rangle = 0$$

$$\langle d^* \alpha, \beta \rangle$$

$$\Rightarrow d^* \alpha = 0.$$

Note:  $\Delta \alpha = 0 \Leftrightarrow d\alpha = d^* \alpha = 0.$

Rmk: • We started with a class  $[\alpha] \in H_{dR}^k(M)$ .

The auxillary structure, (in this case, the metric on  $M$ )  $\mapsto \|\cdot\|$ .

$\mapsto$  unique (harmonic) representative of the class  $[\alpha]$ .

• Simpler example:  $W \subseteq V$  subspace of vector space.

Consider a class  $[\alpha] \in V/W$ .

A priori, there is no preferred choice of representative.

An auxiliary structure, an inner product on  $V \mapsto S(v) = \|v\|$   
 $\mapsto$  unique representative of  $[\alpha]$ , (its orthogonal proj. onto  $W^\perp$ ).

- Donaldson-Uhlenbeck-Yau:

Consider the space of connections on a hol. v.b.  $E \rightarrow X$ .

The auxiliary structure. (a Kähler metric on  $X$ ) gives a preferred choice of connection (Hermitian-Yang-Mills connection), when  $E$  is polystable.

- Thomas-Yau Conjecture: (slag repr.)

Slogan: "metric"  $\mapsto$  preferred representative in each class.  
(harmonic form, orthogonal projection, ...)